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Homer Reid's Solutions to Goldstein Problems: Chapter 8 2 From this we can immediately identify the T matrix and its inverse: $T = \begin{pmatrix} 2k & 2k \\ 2k & 2k \end{pmatrix} a + bq$ $T^{-1} = \begin{pmatrix} a + bq & 2 \\ 1 & 4 - k \end{pmatrix} \begin{pmatrix} a + bq & 2 \\ 1 & 2 \end{pmatrix} a + bq$ $\begin{pmatrix} 2 & 1 \\ -k & 2 \end{pmatrix} \begin{pmatrix} 2 & 1 \\ -k & 2 \end{pmatrix}$ Then the Hamiltonian is $H = \frac{1}{2} a + bq \begin{pmatrix} 2 & 1 \\ 4 - k & 2 \end{pmatrix} p_1 p_2 a + bq \begin{pmatrix} 2 & 1 \\ -k & 2 \end{pmatrix} p_1 p_2 - k \begin{pmatrix} 1 & q \\ 2 & 1 \end{pmatrix} = a + bq \begin{pmatrix} 2 & 1 \\ 4 - k & 2 \end{pmatrix} p_2 \begin{pmatrix} 1 & a + bq \\ 2 & 1 - k \end{pmatrix} p_1 p_2 + p_2^2 - k \begin{pmatrix} 1 & q \\ 2 & 1 \end{pmatrix}$.

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$1 + k \cdot 2)(Q + b \sin(\!t))$ (54) The Hamiltonian is now explicitly dependent on time, and hence is not conserved, as is confirmed by the fact that $dH/dt \neq 0$. The energy is given by $E = T + V = \frac{1}{2} (Q + b \cos(\!t))^2 + \frac{1}{2} (k \cdot 1 + k \cdot 2)(Q + b \sin(\!t))^2$ (55) So, $dE/dt = m(Q + b \cos(\!t))(Q + b \sin(\!t)) + (k \cdot 1 + k \cdot 2)(Q + b \sin(\!t))^2$.

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Solutions Goldstein Chapter 9. CHAPTER 9 – CANONICAL TRANSFORMATIONS

DERIVATIONS: 9.4. Show directly that the transformation is canonical. 9.4. Sol. We are given a transformation as follows, We know that the fundamental Poisson Brackets of the transformed variables have the same value when evaluated with respect to any canonical coordinate set.

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We use the first constraint to solve for the coordinate r : $r = R + a$, $r' = r'' = 0$. We use this solution in Lagrange's equations for r , θ : $m(R + a)r'' + mg \sin \theta = m(R + a)2 r'' + mg(R + a) \cos \theta = \mu(R + a)$
(6) (7) We use the rolling constraint to find an expression for θ as a function of r : $\theta = \frac{r - a}{R + a} + \theta_0$ (8)

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